

# Aspects géométriques et probabilistes des gaz de Coulomb



David García Zelada

28 Juin 2019

# Geometric and probabilistic aspects of Coulomb gases



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# Outline

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A general model and a Laplace principle

Coulomb gas and concentration of measure

Background ensembles and limiting point processes

Selected perspectives

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- $\mathcal{Z}_n = \int_{M^n} \exp(-n\beta_n H_n) d\sigma^{\otimes n} \in (0, \infty)$ , partition function.

## Example: Two-body interaction model

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$G : M \times M \rightarrow (-\infty, \infty]$  bounded from below,

$$H_n(x_1, \dots, x_n) = \frac{1}{n^2} \sum_{i < j} G(x_i, x_j).$$

## Question: macroscopic behavior

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**How  $\{H_n\}_{n \in \mathbb{N}}$  determines the limit of  $\{\hat{\mu}_n\}_{n \in \mathbb{N}}$ ?**

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- whenever  $\frac{1}{n} \sum_{i=1}^n \delta_{x_i} \rightarrow \mu$

$$\liminf_{n \rightarrow \infty} H_n(x_1, \dots, x_n) \geq H(\mu).$$

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If  $G$  is lower semicontinuous,

$$H(\mu) = \frac{1}{2} \int_{M \times M} G(x, y) d\mu(x) d\mu(y)$$

is the macroscopic limit of  $\{H_n\}_{n \in \mathbb{N}}$ .

## Relative entropy and free energy

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### Definition (Relative entropy)

Let  $\mu \in \mathcal{P}(M)$ . If  $d\mu = \rho d\sigma$

$$D(\mu \parallel \sigma) = \int_M \rho \log \rho d\sigma.$$

Otherwise,  $D(\mu \parallel \sigma) = \infty$ .

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### Definition (Free energy)

Let  $\beta \in (0, \infty)$  and  $H : \mathcal{P}(M) \rightarrow (-\infty, \infty]$ . For  $\mu \in \mathcal{P}(M)$

$$F_\beta(\mu) = H(\mu) + \frac{1}{\beta} D(\mu\|\sigma).$$

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$$\lim_{n \rightarrow \infty} \frac{1}{n\beta_n} \log \mathbb{E} \left[ e^{-n\beta_n f(\hat{\mu}_n)} \right] = - \inf_{\mu \in \mathcal{P}(M)} \{f(\mu) + I_\beta(\mu)\}.$$

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*History: [Sanov '57],*

*[Messer & Spohn '82], [Caglioti, Lions, Marchioro & Pulvirenti '92],*

*[Bodineau & Guionnet '99], [Dupuis, Laschos & Ramanan '15] and [Berman '18].*

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- $H$  **regular**: If  $H(\mu) < \infty$

$$\exists \{\mu_n\}_{n \in \mathbb{N}} \text{ s.t. } \mu_n \rightarrow \mu,$$

$$\forall n, D(\mu_n \| \sigma) < \infty \text{ and } \lim_{n \rightarrow \infty} H(\mu_n) = H(\mu).$$

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If  $G(x, y) \rightarrow \infty$  whenever  $x, y \rightarrow \infty$  at the same time then

$\{H_n\}_{n \in \mathbb{N}}$  is confining.

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*History: [Ben Arous & Guionnet '97], [Ben Arous & Zeitouni '98], [Hiai & Petz '98], [Zeitouni & Zelditch '10], [Hardy '12], [Chafaï Gozlan & Zitt '14], [Dupuis, Laschos & Ramanan '15] and [Berman '18].*

## Idea of the proof

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*Motivation: Weak convergence approach to large deviations of Dupuis and Ellis.*

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**Use macroscopic convergence and properties of the entropy!**

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Let  $\Lambda$  be a differentiable signed measure on  $M$  s.t.  $\Lambda(M) = 1$ .  
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Then

$$d\mathbb{P}_n(z_1, \dots, z_n) = \frac{1}{Z_n} \prod_{i < j} |z_i - z_j|^{\frac{\beta n}{2\pi n}} e^{-\left(\frac{\beta n}{2\pi n}(n-1)+4\right) \sum_{i=1}^n V(z_i)} d\ell_{\mathbb{C}^n}(z_1, \dots, z_n).$$

## Convergence of empirical measures

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If  $\beta_n \rightarrow \infty$  and  $(X_1, \dots, X_n) \sim \mathbb{P}_n$  then

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Definition (Kantorovich-Wasserstein distance)

$$W_1 : \mathcal{P}(M) \times \mathcal{P}(M) \rightarrow [0, \infty)$$

$$W_1(\mu, \nu) = \inf_{X \sim \mu, Y \sim \nu} \mathbb{E} [d(X, Y)].$$



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This distance metrizes the weak convergence!

$$W_1(\hat{\mu}_n, \mu_{\text{eq}}) \xrightarrow{a.s.} 0.$$

## Question: concentration of measure

---

LDP implies  $\exists \mathcal{I}_r > 0$  and  $\exists \varepsilon_n$  s.t.  $\varepsilon_n \rightarrow 0$

$$\mathbb{P} (W_1(\hat{\mu}_n, \mu_{\text{eq}}) \geq r) \leq \exp (-n\beta_n \mathcal{I}_r + n\beta_n \varepsilon_n) .$$

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**Can we find a simple  $\mathcal{I}_r$  and  $\varepsilon_n$ ?**

## Concentration for Coulomb gases

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### Theorem (Concentration inequality [GZ])

*If  $\dim M = 2$  then  $\exists C = C(M, \Lambda)$  s.t.*

$$\mathbb{P}(W_1(\hat{\mu}_n, \mu_{\text{eq}}) \geq r) \leq \exp \left( -n\beta_n \frac{r^2}{4} + \frac{\beta_n}{8\pi} \log(n) + nD(\mu_{\text{eq}} \parallel \sigma) + C\beta_n \right).$$

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If  $\dim M \geq 3$  then  $\exists C = C(M, \Lambda)$  s.t.

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*History: [Maïda & Maurel-Segala '12], [Rougerie & Serfaty '16], [Chafaï, Hardy & Maïda '18] and [Berman '19].*

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$\exists R_n : M^n \rightarrow \mathcal{P}(M)$  and  $C > 0$  s.t.

$$H_n(x_1, \dots, x_n) \geq H(R_n(x_1, \dots, x_n)) - \frac{1}{8\pi n} \log(n) - \frac{C}{n} \quad \text{and}$$

$$W_1 \left( R_n(x_1, \dots, x_n), \frac{1}{n} \sum_{i=1}^n \delta_{x_i} \right) \leq \frac{C}{\sqrt{n}}.$$

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Let  $p_t(x, y)$  be the heat kernel and

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## Outline

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A general model and a Laplace principle

Coulomb gas and concentration of measure

Background ensembles and limiting point processes

Selected perspectives

## Determinantal case

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Let  $\nu \in \mathcal{P}(\mathbb{C})$  be rotationally invariant so that

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If  $(X_1, \dots, X_n) \sim \mathbb{P}_n$  and  $f : \mathbb{C} \rightarrow \mathbb{R}$  is bounded continuous then

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## Question: uncharged regions

---

If  $\text{supp } f \cap \text{supp } \nu = \emptyset$  then

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**What is**  $\lim_{n \rightarrow \infty} \sum_{i=1}^n f(X_i)$ ?

## Possible cases

---

Connected components of  $\mathbb{C} \setminus \text{supp } \nu$ :

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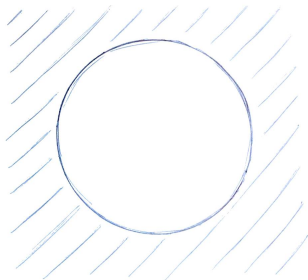
*A family of possibilities (up to rescaling)*

- $\mathbb{D} = D_1$ ,
- $\mathbb{C} \setminus \bar{\mathbb{D}}$  and
- $\mathbb{A}_R = D_R \setminus \bar{\mathbb{D}}$  for  $R > 1$ .



$\mathbb{D}$  is a connected component

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Theorem (Limiting Bergman point process [Butez & GZ])

$\forall f : \mathbb{C} \rightarrow \mathbb{R}$  *continuous and s.t.*  $\text{supp } f \subset \mathbb{D}$

$$\lim_{n \rightarrow \infty} \sum_{i=1}^n f(X_i) = \sum_{z \in \mathcal{B}_{\mathbb{D}}} f(z)$$

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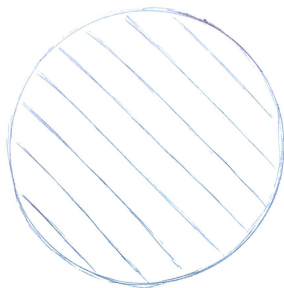
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*History: [Ameur, Kang & Makarov '18].*

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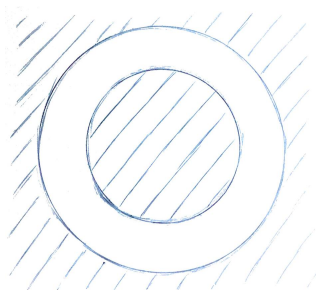
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## $\mathbb{A}_R$ is a connected component

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### Theorem (Limiting Bergman point processes on $\mathbb{A}_R$ [GZ])

*Let  $q = \mu(\bar{\mathbb{D}})$  and take a subsequence of integers s.t.*

$$\exp(2\pi(n+1)qi) \rightarrow \exp(2\pi\gamma i)$$

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in law where  $\mathcal{B}_{\mathbb{A}_R, \gamma}$  is a DPP on  $\mathbb{A}_R$  associated to the kernel of

**the orthogonal projection onto the space of holomorphic functions of  $\mathbb{A}_R$**

with weight  $|z|^{-2\gamma}$ .

## Idea of the proof

---

$\sum_{i=1}^n \delta_{X_i}$  is a DPP with kernel

$$K_n(z, w) = \sum_{l=0}^{n-1} a_l z^l \bar{w}^l e^{-(n+1)V(z)} e^{-(n+1)V(w)}$$

where

$$a_l = \left( \int_{\mathbb{C}} |z|^{2l} e^{-2(n+1)V(z)} d\ell_{\mathbb{C}}(z) \right)^{-1}.$$

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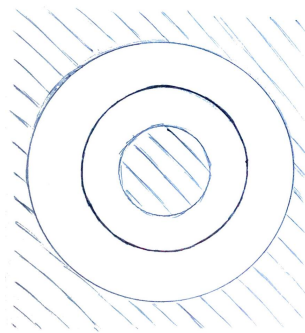
For the annulus,

$$\sum_{k=-\lfloor nq \rfloor}^{n-1-\lfloor nq \rfloor} a_{\lfloor nq \rfloor + k} z^k \bar{w}^k |z|^{\lfloor nq \rfloor} |w|^{\lfloor nq \rfloor} e^{-(n+1)V(z)} e^{-(n+1)V(w)}.$$

## Behavior between two uncharged regions

---

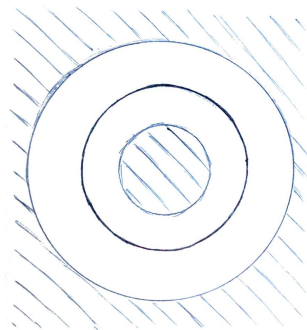
Suppose there are  $R_1 < 1 < R_2$  s.t.  $\nu(R_1 < |z| < R_2) = \nu(\partial\mathbb{D}) > 0$



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Theorem (Behavior at the unit circle [GZ])

$$\lim_{n \rightarrow \infty} \sum_{i=1}^n \delta_{n(X_i-1)} = \xi_{q,Q}$$

where  $\xi_{q,Q}$  is a DPP with kernel

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History: [Sinclair & Yattselev '12] and [Hedenmalm & Wennman '17].

## Idea of the proof

---

Study

$$\rho_n(r) = \frac{1}{n^2} \sum_{k=0}^{n-1} a_k \left(1 + \frac{r}{n}\right)^{\left(\frac{k}{n}\right)n}$$

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$$a_k = \left( \int_{\mathbb{C}} |z|^{2k} e^{-2(n+1)V(z)} d\ell_{\mathbb{C}}(z) \right)^{-1}.$$

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$$\rho_n(r) = \int_0^1 \frac{a_{t^*n}}{n} \left(1 + \frac{r}{n}\right)^{t^*n} dt$$

and take the limit!

## Behavior of the maxima

---

### Theorem (Compactly supported background [Butez & GZ])

If  $\nu(\mathbb{C} \setminus D_R) = 0$  and  $\partial D_R \subset \text{supp } \nu$ ,

$$\forall r \geq R, \quad \lim_{n \rightarrow \infty} \mathbb{P}(\max\{|X_1|, \dots, |X_n|\} < r) = \prod_{k=1}^{\infty} (1 - R^{2k} r^{-2k}).$$

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### Theorem (Non-compactly supported background [Butez & GZ])

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*History: [Rider '03], [Chafaï & Peche '14] and [Jiang & Qi '17].*

# Outline

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A general model and a Laplace principle

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## Macroscopic behavior

---

Recall the Coulomb gas law

$$d\mathbb{P}_n(z_1, \dots, z_n) = \frac{1}{\mathcal{Z}_n} \prod_{i < j} |z_i - z_j|^{\frac{\beta_n}{2\pi n}} e^{-\left(\frac{\beta_n}{2\pi n}(n-1)+4\right) \sum_{i=1}^n V(z_i)} d\ell_{\mathbb{C}^n}(z_1, \dots, z_n).$$

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**It makes sense for any  $\beta_n > -8\pi$ !**

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---

Recall the Coulomb gas law

$$d\mathbb{P}_n(z_1, \dots, z_n) = \frac{1}{\mathcal{Z}_n} \prod_{i < j} |z_i - z_j|^{\frac{\beta_n}{2\pi n}} e^{-\left(\frac{\beta_n}{2\pi n}(n-1)+4\right) \sum_{i=1}^n V(z_i)} d\ell_{\mathbb{C}^n}(z_1, \dots, z_n).$$

**It makes sense for any  $\beta_n > -8\pi$ !**

What happens when  $\beta_n \downarrow -8\pi$ ?

## Fluctuations

---

Suppose the background measure on  $M$  is  $\Lambda = \sigma$  so that if

$$(X_1, \dots, X_n) \sim \text{Coulomb gas on } M$$

then

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=1}^n \delta_{X_i} = \sigma.$$

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Suppose  $\int_M f d\sigma = 0$  and  $\beta_n = n\beta$ .

$$\text{What is } \lim_{n \rightarrow \infty} \sum_{i=1}^n f(X_i)?$$

## Uncharged region

---

Suppose  $\text{supp } f \cap \text{supp } \nu = \emptyset$ . We know

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- What happens in the non-determinantal case?

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*Three open questions:*

- What happens if  $\nu$  is not radial?
- What happens in the non-determinantal case?
- What happens if we only assume  $f|_{\text{supp } \nu} = 0$ ?

Thank you for your attention!

Merci pour votre attention!

¡Gracias por su atención!

*Extra frames*

## Proof of the annulus case

## Proof Annulus case, p.1

---

$\sum_{i=1}^n \delta_{X_i}$  is a DPP with kernel

$$K_n(z, w) = \sum_{l=0}^{n-1} a_l z^l \bar{w}^l e^{-(n+1)V(z)} e^{-(n+1)V(w)}$$

where

$$a_l = \left( \int_{\mathbb{C}} |z|^{2l} e^{-2(n+1)V(z)} d\ell_{\mathbb{C}}(z) \right)^{-1}.$$

## Proof Annulus case, p.2

---

Since  $q = \mu(\mathbb{D})$ , we have

$$V(z) = q \log |z| \quad \text{for } z \in \bar{\mathbb{A}}_R$$

and

$$V(z) > q \log |z| \quad \text{if } z \notin \bar{\mathbb{A}}_R.$$

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We assume  $\lim_{n \rightarrow \infty} \{(n+1)q - \lfloor nq \rfloor\} = \gamma$  so that

$$\lim_{n \rightarrow \infty} |z|^{\lfloor nq \rfloor} e^{-(n+1)V(z)} = |z|^{-\gamma}.$$

## Proof Annulus case, p.3

---

By a change of index

$$\begin{aligned} \left(\frac{|z|}{z}\right)^{\lfloor nq \rfloor} K_n(z, w) \left(\frac{|w|}{\bar{w}}\right)^{\lfloor nq \rfloor} \\ = \sum_{k=-\lfloor nq \rfloor}^{n-1-\lfloor nq \rfloor} a_{\lfloor nq \rfloor+k} z^k \bar{w}^k |z|^{\lfloor nq \rfloor} |w|^{\lfloor nq \rfloor} e^{-(n+1)V(z)} e^{-(n+1)V(w)}. \end{aligned}$$

## Proof Annulus case, p.3

---

By a change of index

$$\begin{aligned} \left(\frac{|z|}{z}\right)^{[nq]} K_n(z, w) \left(\frac{|w|}{\bar{w}}\right)^{[nq]} \\ = \sum_{k=-[nq]}^{n-1-[nq]} a_{[nq]+k} z^k \bar{w}^k |z|^{[nq]} |w|^{[nq]} e^{-(n+1)V(z)} e^{-(n+1)V(w)}. \end{aligned}$$

$$\lim_{n \rightarrow \infty} \left(\frac{|z|}{z}\right)^{[nq]} K_n(z, w) \left(\frac{|w|}{\bar{w}}\right)^{[nq]} = \sum_{k \in \mathbb{Z}} b_k z^k \bar{w}^k |z|^{-\gamma} |w|^{-\gamma}$$

where

$$b_k = \lim_{n \rightarrow \infty} a_{[nq]+k} = \left( \int_{\mathbb{A}_R} |z|^{2k} |z|^{-2\gamma} d\ell_{\mathbb{C}}(z) \right)^{-1}.$$

Proof of the behavior at the unit circle

## Proof of the behavior at the unit circle, p.1

---

Study

$$\rho_n(r) = \frac{1}{n^2} \sum_{k=0}^{n-1} a_k \left(1 + \frac{r}{n}\right)^{\left(\frac{k}{n}\right)n}$$

where

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Define  $t^*(t) = \lfloor tn \rfloor / n$ .

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Define  $t^*(t) = \lfloor tn \rfloor / n$ . Then

$$\rho(r) = \int_0^1 \frac{a_{t^*n}}{n} \left(1 + \frac{r}{n}\right)^{t^*n} dt.$$

## Proof of the behavior at the unit circle, p.2

---

Since

$$\lim_{n \rightarrow \infty} \frac{a_{t^*n}}{n} = \frac{1}{\pi} \frac{(Q-t)(t-q)}{Q-q} 1_{(q,Q)}(t)$$



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we have

$$\lim_{n \rightarrow \infty} \rho_n(r) = \frac{1}{\pi} \int_q^Q \frac{(Q-t)(t-q)}{Q-q} e^{rt} dt.$$

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By Montel's theorem, the same for  $r \in \mathbb{C}$ . Then, choose

$$r_n = z + \bar{w} + \frac{z\bar{w}}{n}$$

that satisfies

$$\lim_{n \rightarrow \infty} r_n = z + \bar{w} \quad \text{uniformly on compact sets.}$$